



# TIME-DEPENDENT DYNAMICAL ENERGY ANALYSIS VIA CONVOLUTION QUADRATURE

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## ABSTRACT

Dynamical Energy Analysis (DEA) was introduced in 2009 as a novel method for predicting high-frequency acoustic and vibrational energy distributions. The Convolution Quadrature (CQ) method provides a link between the frequency domain and fully time-dependent solutions by means of the  $Z$ -transform when the time-dependent solution can be expressed as a convolution in time. In this work we detail how DEA can be formulated in the time domain by means of a convolution integral operator and apply the CQ method to discretise in time. The space and momentum variables may be approximated using the same approaches that have previously been implemented in frequency domain DEA. The final result is a fully time-dependent DEA method that can track the propagation of high-frequency transient signals through phase-space.

**Keywords:** *Dynamical Energy Analysis, Boundary Integral Method, Convolution Quadrature, High-Frequency, Ray-tracing*

## 1. INTRODUCTION

Boundary integral methods for modelling time-dependent wave propagation were originally proposed in the 1960s [1, 2]. The considerable increase in available computer power during the latter part of the twentieth century made numerical solutions over longer time intervals feasible, and with this advance long-time instabilities in the numerical solutions also became evident [3–5]. The cause

of these instabilities has been linked to internal resonances of the wave scatterer for exterior problems [4], or the region being modelled for interior problems. For this reason, combined field integral equations, such as the time-dependent Burton and Miller formulation, have been proposed to tackle these stability issues [6, 7]. However, these formulations introduce additional computational overheads and the need to evaluate hypersingular boundary integral operators. An alternative is to apply the CQ method, see for example Refs. [8–10], which is able to provide stable results based on standard integral equation formulations. The reason for the preferable stability properties of CQ essentially relate to the reconstruction of the time domain solution, or alternatively the time domain boundary integral operator, through a numerical inverse Laplace transform where the contour is taken over Laplace domain frequencies with strictly positive real part. Since the resonances lie on the imaginary axis in the Laplace domain, then they do not effect the result in the time domain.

For high-frequency time-dependent wave problems, such as those arising in the interior acoustics of large spaces like concert halls, ray-tracing methods are often preferred to full wave models, see for example [11]. Traditional ray based methods work well for applications where only a few reflections need to be considered, but not so well for problems including multiple scattering and chaotic dynamics. In this case, multiple reflections of the rays can give an exponentially growing number of trajectories to track. Dynamical Energy Analysis (DEA) is a phase-space boundary integral method that models wave energy densities [12]. DEA is a frequency domain method formed by seeking solutions of the stationary Liouville equation, circumventing issues regarding the exponentially growing number of rays to track as time increases [13].

Time-domain simulations are important for various applications such as predicting radar cross sections,

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shock-responses and, of course, auralisation. In this proceedings paper we outline a methodology for extending DEA to the time-domain based on the CQ method. The computational cost of time-domain DEA should scale only linearly with the modelled time period regardless of the ray dynamics, comparing favourably with conventional ray-tracers.

## 2. OUTLINE OF METHODOLOGY

In order to develop a time dependent DEA method for a domain  $\Omega$  containing an acoustic medium with speed of sound  $c$ , the first step is to reformulate the DEA phase-space boundary integral operator to have explicit time dependence. The result is a one-sided convolution (in time  $t$ ) operator  $\mathcal{B}$  given by

$$(\mathcal{B}\rho_0)(t, s, p) = \int_{c|p'| \leq 1} \int_{\Gamma} (k * \rho_0)(t, s', p') d\Gamma(s') dp', \quad (1)$$

which is applied to a specified initial density distribution of rays  $\rho_0$  on the boundary  $\Gamma$ . Here  $k$  is the time-dependent kernel of our boundary integral operator, which is given by a multidimensional Dirac delta generalised function specifying the propagation of a ray through time, position and momentum. The variables  $(s', p')$  relate respectively to the position and momentum of the starting position of a ray emanating from  $\Gamma$  and  $(s, p)$  correspond to the arrival position and momentum on  $\Gamma$ , respectively, following a specular reflection. Note that a damping factor must be applied to obtain convergence in frequency domain DEA, but this is not necessary in the time-dependent formulation owing to the fact that we only model a finite time duration  $[0, T]$ .

The CQ method can be applied for the time discretisation of one-sided convolution operators such as  $\mathcal{B}$  (1), see for example [8–10]. After applying the time-discretisation, to obtain a fully discrete problem we then need to discretise in the position and momentum variables. Here we may use any of the discretisation methods previously applied for frequency domain DEA amongst others, see for example [12–14].

## 3. CONCLUSIONS AND FUTURE WORK

In this short paper we have motivated and outlined a methodology to extend the DEA method for time-dependent problems. We will present numerical results based on this work at the conference. For these results we will apply the position and momentum discretisation

methods from Ref. [14], since this allows for a verification of the time discretisation method in simple examples for which we can derive an exact solution. We will use these examples to study whether the order of convergence in time is consistent with the expected behaviour of the applied time-stepping approach. We will also present examples for which this choice of position and momentum discretisation is less favourable and it may be preferable to use the methods in Ref. [13], for example, instead.

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